

Is ocean surface wind stress key in the long term predictability of the atmosphere?

Stéphane Vannitsem

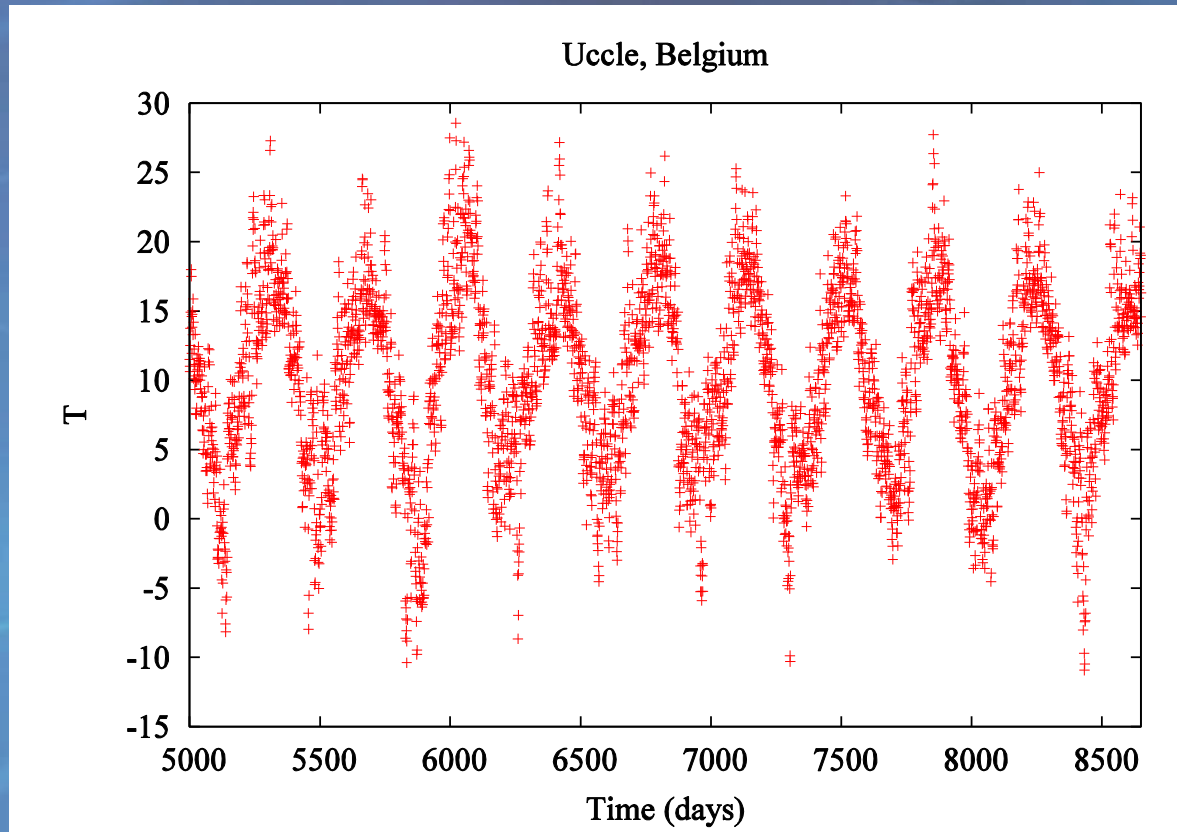
Royal Meteorological Institute of Belgium

Paris, October 9, 2018

Collaborations: L. De Cruz, J. Demaeyer, M. Ghil, V. Lucarini,
S. Schubert, R. Solé-Pomies

Introduction

Weather variability



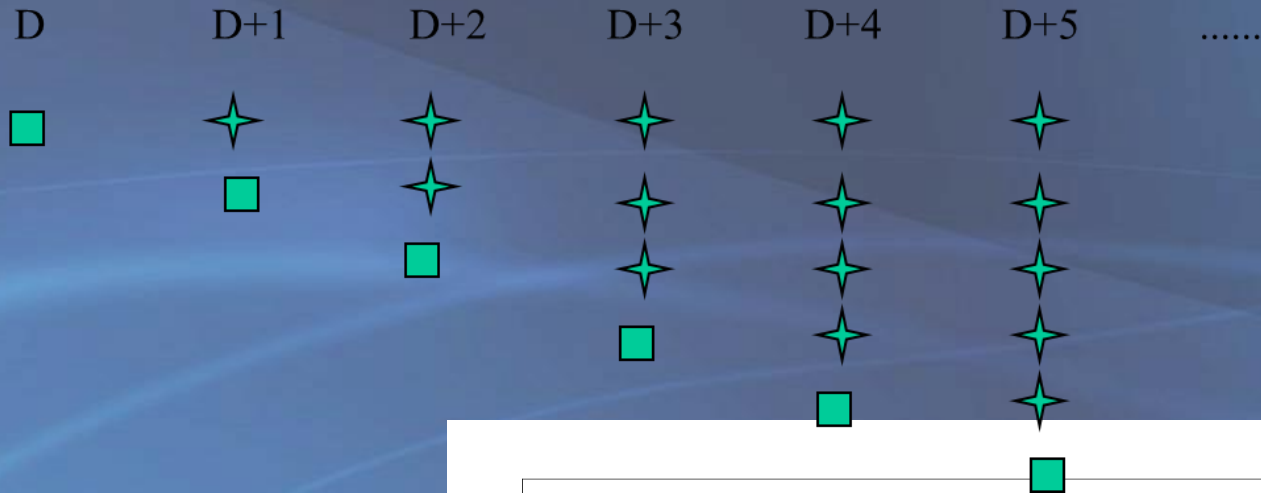
Predictability

The property of sensitivity to initial (and model) uncertainties at the origin of the degradation of the quality of forecasts of atmospheric flows

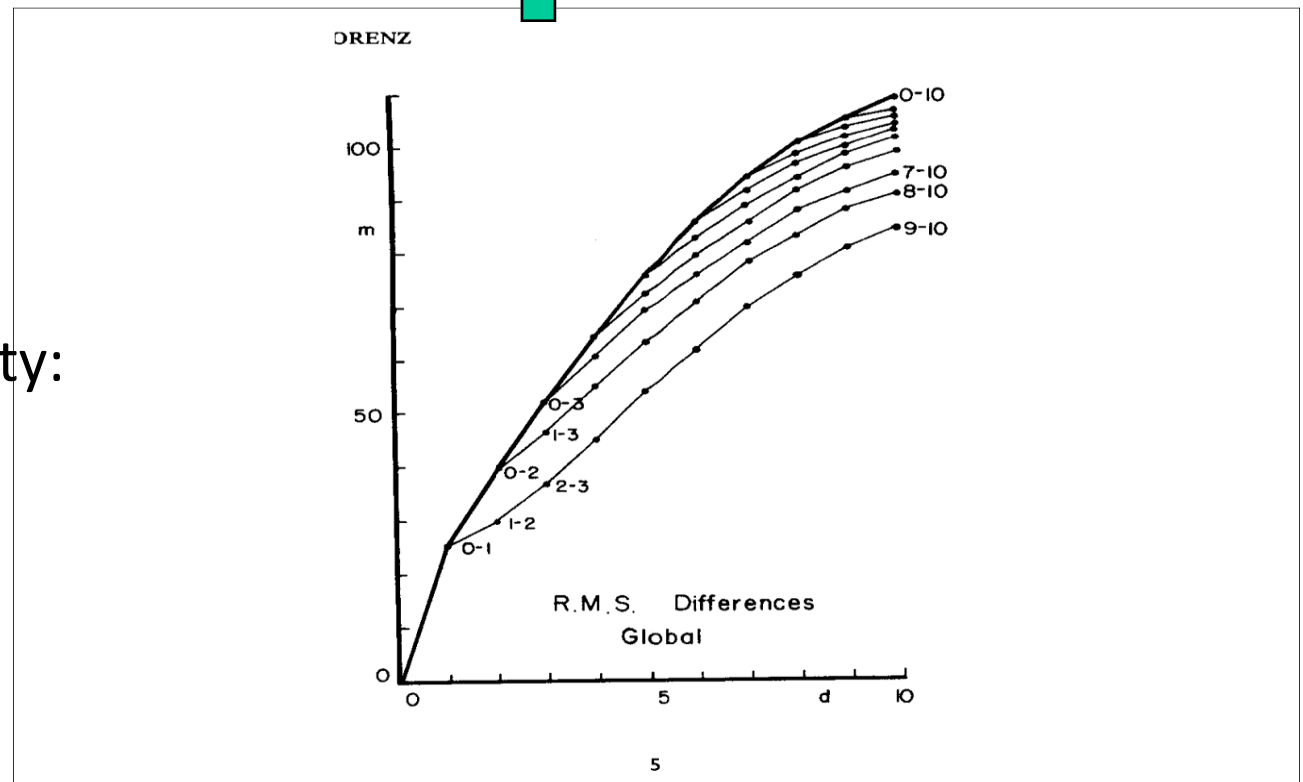
Property already recognized by
Thompson (1957, *Tellus*, 9) and Lorenz (1963)

From a mathematical point of view: Poincaré (1888; 1908, *Science et méthode*)



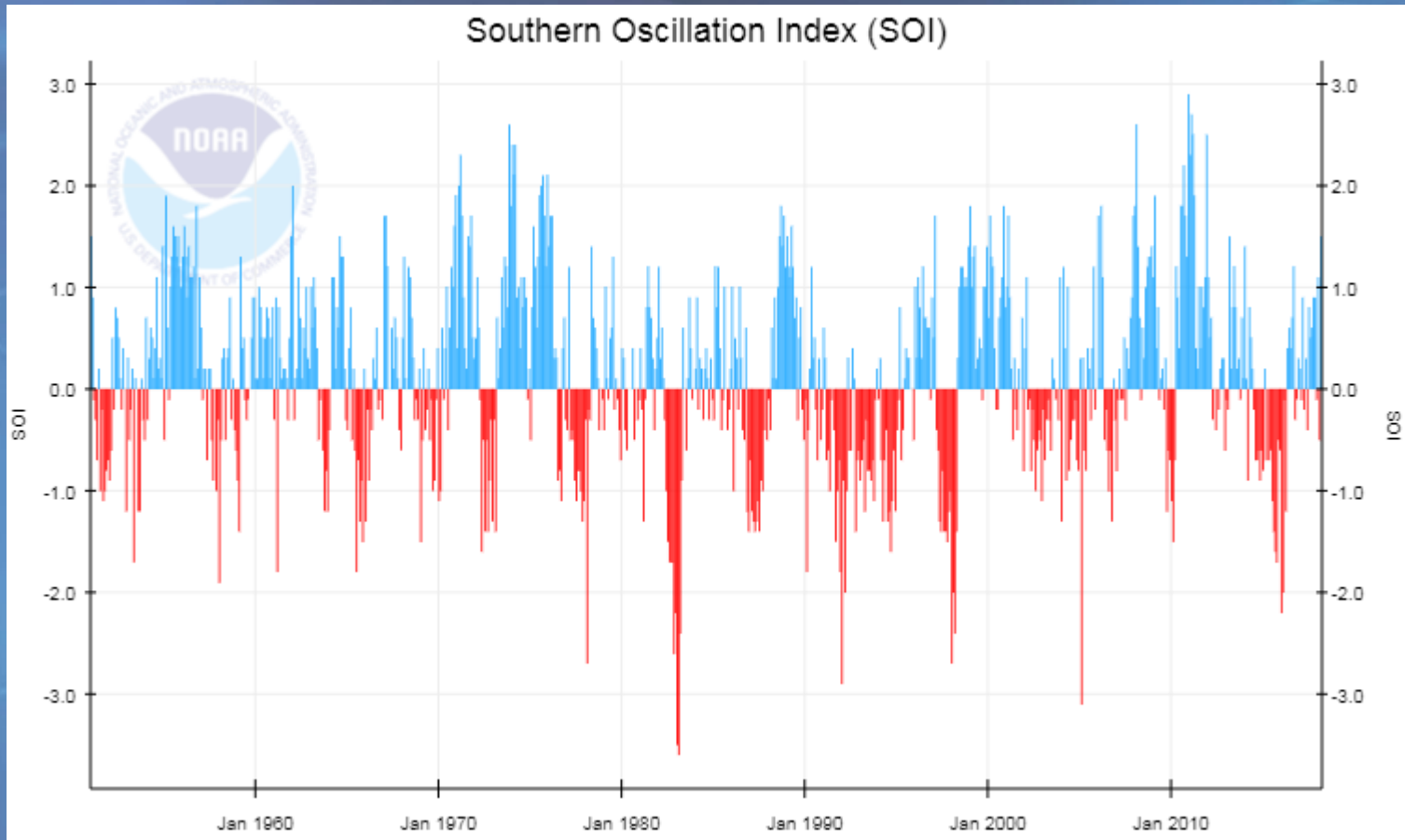


Limited predictability:
about 10-15 days



Climate variability and predictability?

One important signal: Southern Oscillation Index

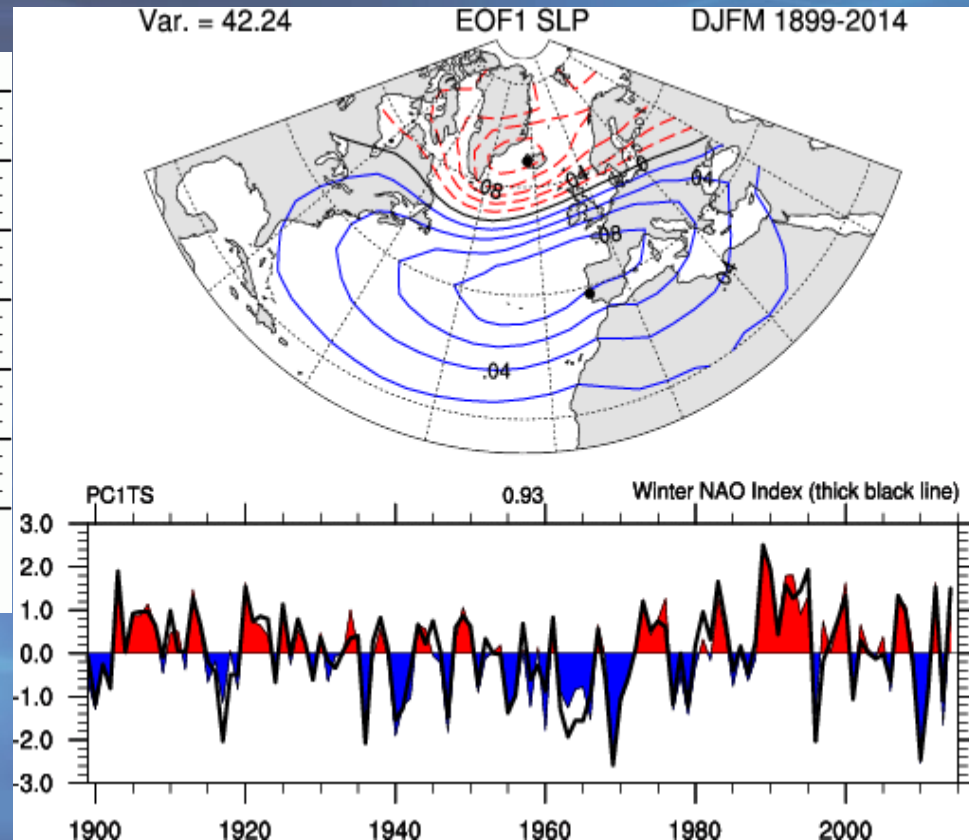
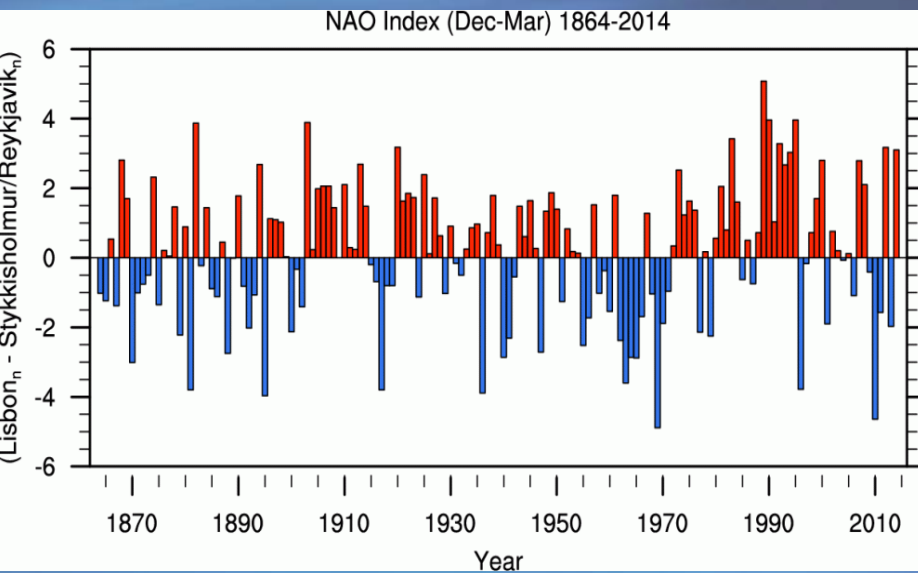


Associated with the development of El-Nino and La-Nina

El-Nino-Southern-Oscillation (ENSO)

Climate variability and predictability?

North Atlantic Oscillation (NAO)



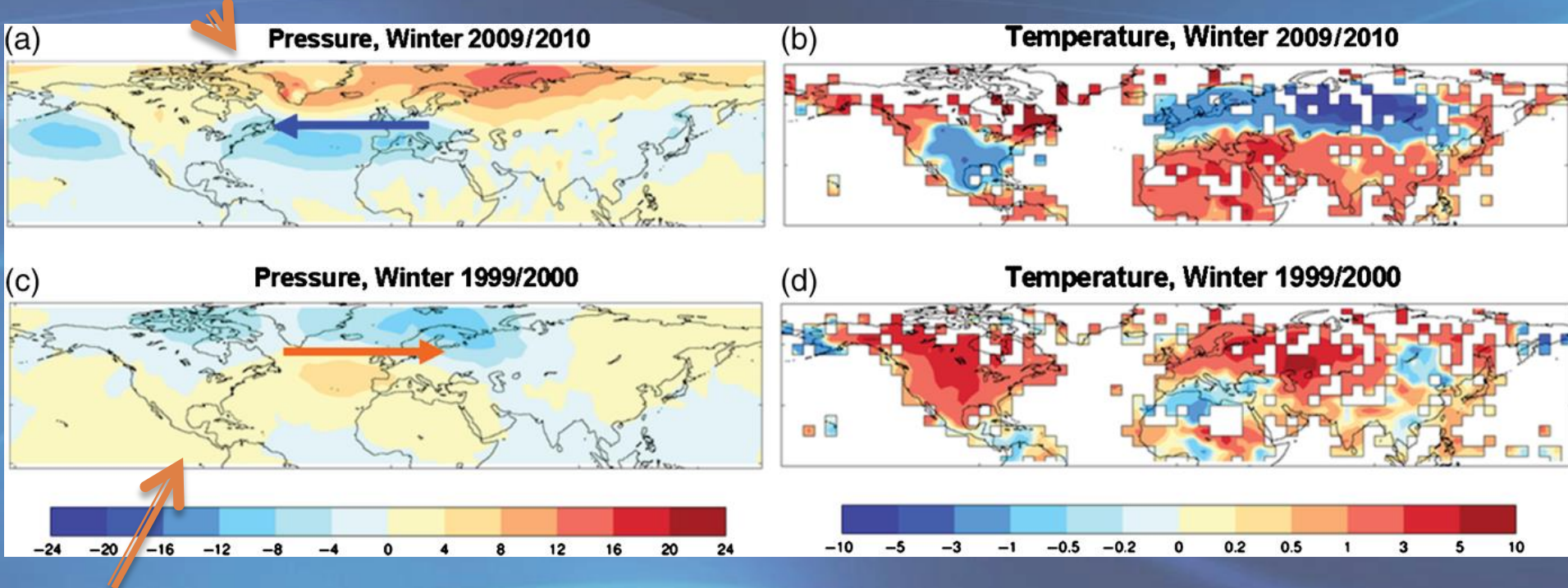
Positive NAO: Larger pressure difference between Lisbon and Reykjavik

Negative NAO: Smaller pressure difference

Climate variability and predictability?

Implications?

Negative NAO



Positive NAO

- EMBEDDED IN THE NORTH ATLANTIC OCEAN DYNAMICS?
- TELECONNECTION WITH THE TROPICAL PACIFIC?
- OTHER PROCESSES?

Smith et al, 2014

General objective

- To characterize the predictability of the atmosphere on seasonal, inter-annual and decadal time scales

Strategy

- Development of reduced-order climate models, and in particular coupled ocean-atmosphere models
- Analysis of the predictability of atmospheric and climate models of various resolutions

Aim of the presentation

Analyze the mechanisms at the origin of the development of long term (potential) predictability of the atmosphere

To this aim, we are using a coupled reduced-order ocean-atmosphere model, under the assumption that the ocean is playing a crucial role on the development of low-frequency variability and long term predictability

More specifically we will use:

- **Version of the MAOOAM model over the Atlantic, called VDDG which was developed in Vannitsem, Demaeyer, De Cruz and Ghil, 2015, Physica D.**
- **New version of the coupled ocean-atmosphere model, called MAO(S)OAM with different boundary conditions**

Low-order modelling

These are simplified models containing:

- **Key ingredients** of the system's dynamics
- Developed on an **appropriate basis** at the scale of interest

ex: Use of empirical orthogonal functions
 Truncated Fourier series

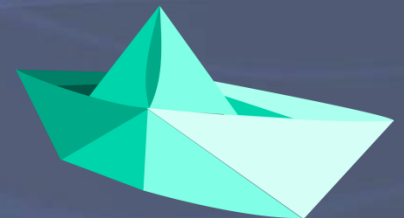
Many models were developed (see e.g. Sprott, 2010)

Procedure

$$\varphi = \sum_i^N C_i F_i$$

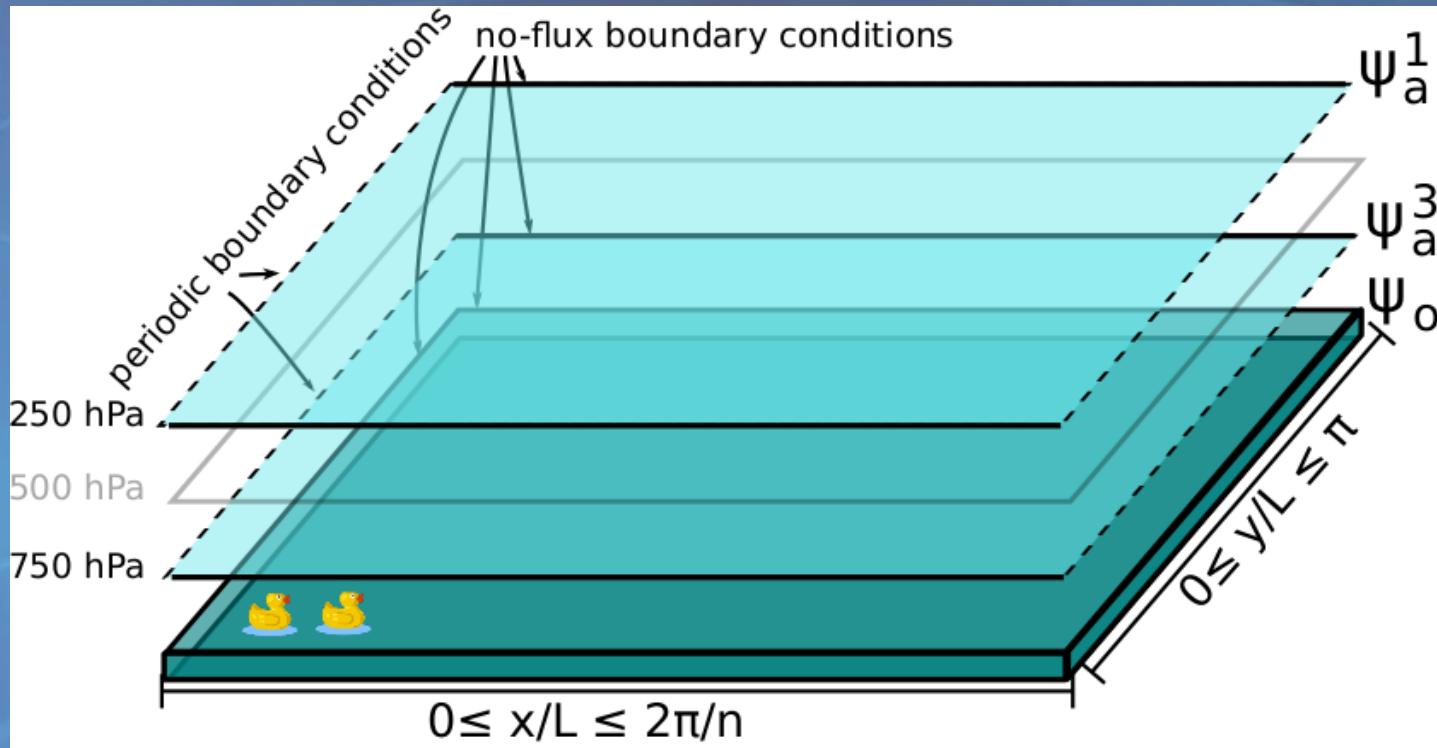
$$(f, g) = \iint_{\mathcal{V}} dx dy f^* g$$

$$\left(F_i, \frac{\partial \varphi}{\partial t} \right) = (F_i, G(\varphi, \nabla \varphi, \dots))$$



An idealized low-order coupled ocean-atmosphere model

- QG model for both the ocean and the atmosphere



Vannitsem et al, 2015, Physica D, 309, 71-85, 2015, (VDDG)

De Cruz et al 2016, Geosci. Model Develop, 9, 2793-2808, 2016. (MAOOAM)

Latitudinal dependence
of the radiative input

$$R_0 + S_0 \sqrt{2} \cos \gamma$$

Surface friction strength

$$\delta = \frac{d}{f_0} = \frac{C}{\rho H f_0}$$

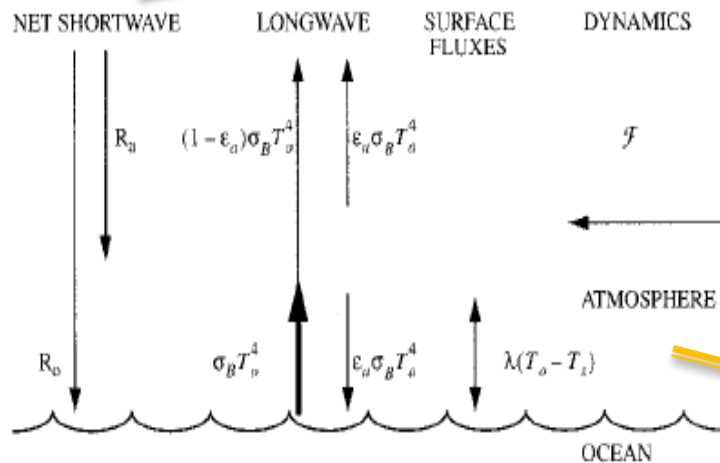
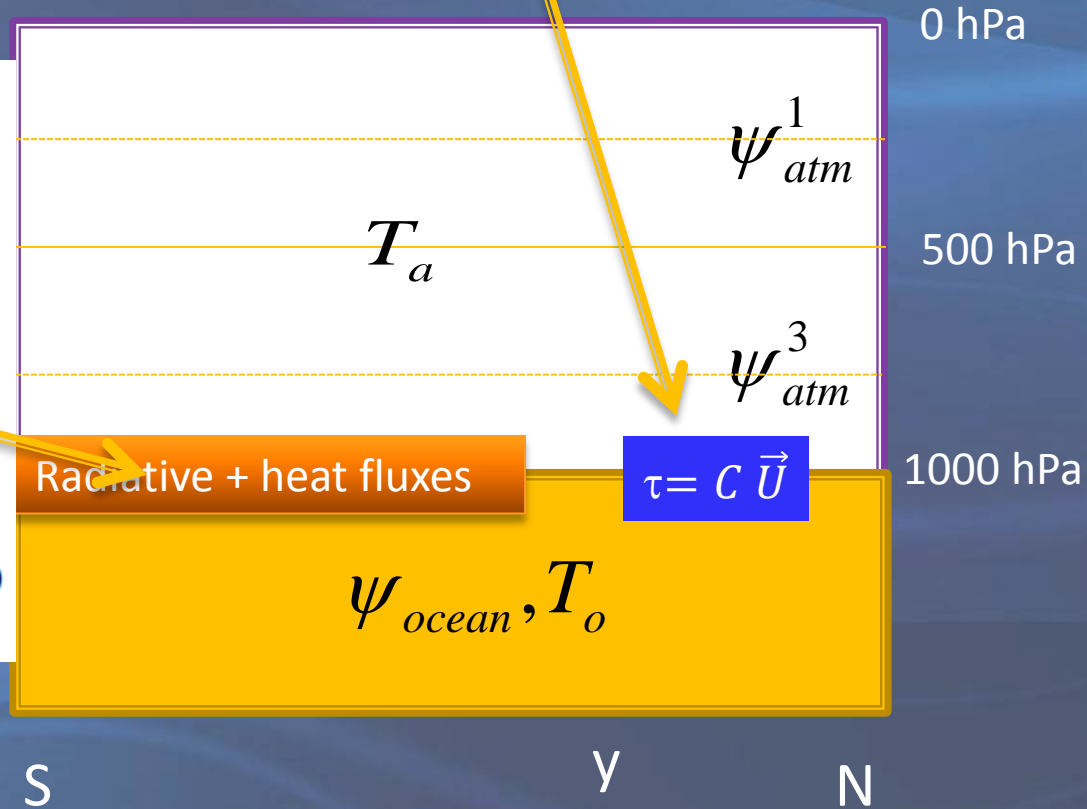


FIG. 2. Diagram of simple energy balance model on which Eqs. (1) and (2) are based. See appendix A for definition of symbols.

Barsugli & Battisti, 1998, JAS



The dynamical equations for the ocean-atmosphere model

For the atmosphere

$$\begin{aligned}\frac{\partial}{\partial t} (\nabla^2 \psi_a^1) + J(\psi_a^1, \nabla^2 \psi_a^1) + \beta \frac{\partial \psi_a^1}{\partial x} &= -k'_d \nabla^2 (\psi^1 - \psi^3) + \frac{f_0}{\Delta p} \omega \\ \frac{\partial}{\partial t} (\nabla^2 \psi_a^3) + J(\psi_a^3, \nabla^2 \psi_a^3) + \beta \frac{\partial \psi_a^3}{\partial x} &= +k'_d \nabla^2 (\psi_a^1 - \psi_a^3) - \frac{f_0}{\Delta p} \omega \\ &\quad - k_d \nabla^2 (\psi_a^3 - \psi_o)\end{aligned}$$

$$\gamma_a \left(\frac{\partial T_a}{\partial t} + J(\psi_a, T_a) - \sigma \omega \frac{p}{R} \right) = -\lambda(T_a - T_o) + E_{a,R}$$

Friction on a moving surface

$$E_{a,R} = \epsilon_a \sigma_B T_o^4 - 2\epsilon_a \sigma_B T_a^4 + R_a$$

For the ocean

$$\frac{\partial}{\partial t} \left(\nabla^2 \psi_o - \frac{\psi_o}{L_R^2} \right) + J(\psi_o, \nabla^2 \psi_o) + \beta \frac{\partial \psi_o}{\partial x} = -r \nabla^2 \psi_o + \frac{\text{curl}_z \tau}{\rho h}$$

$$\gamma_o \left(\frac{\partial T_o}{\partial t} + J(\psi_o, T_o) \right) = -\lambda(T_o - T_a) + E_R$$

Curl of wind stress

$$E_R = -\sigma_B T_o^4 + \epsilon_a \sigma_B T_a^4 + R_o$$

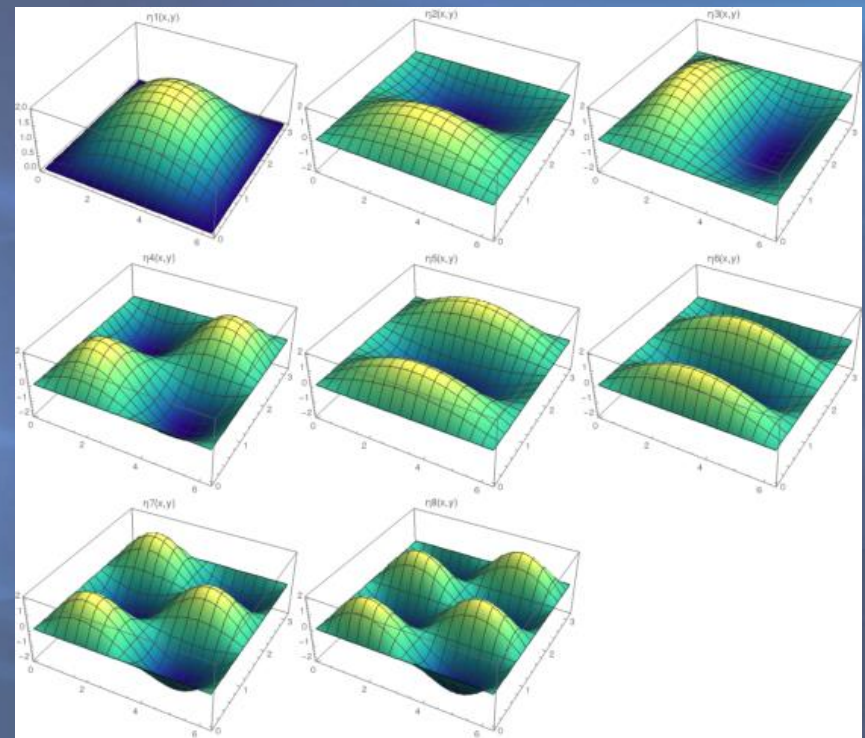
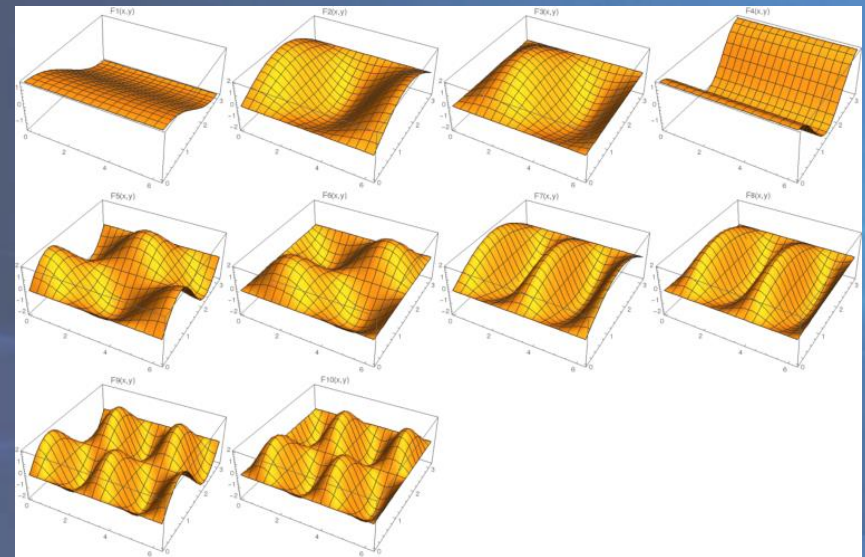
Truncation of Fourier series

$$\psi = \sum_{k=1}^K \psi_k F_k$$

$$\theta = \sum_{k=1}^K \theta_k F_k$$

$$\psi_o = \sum_{i=1}^8 \psi_{o,i} \phi_i, \quad \delta T_o = \sum_{i=1}^8 T_{o,i} \phi_i$$

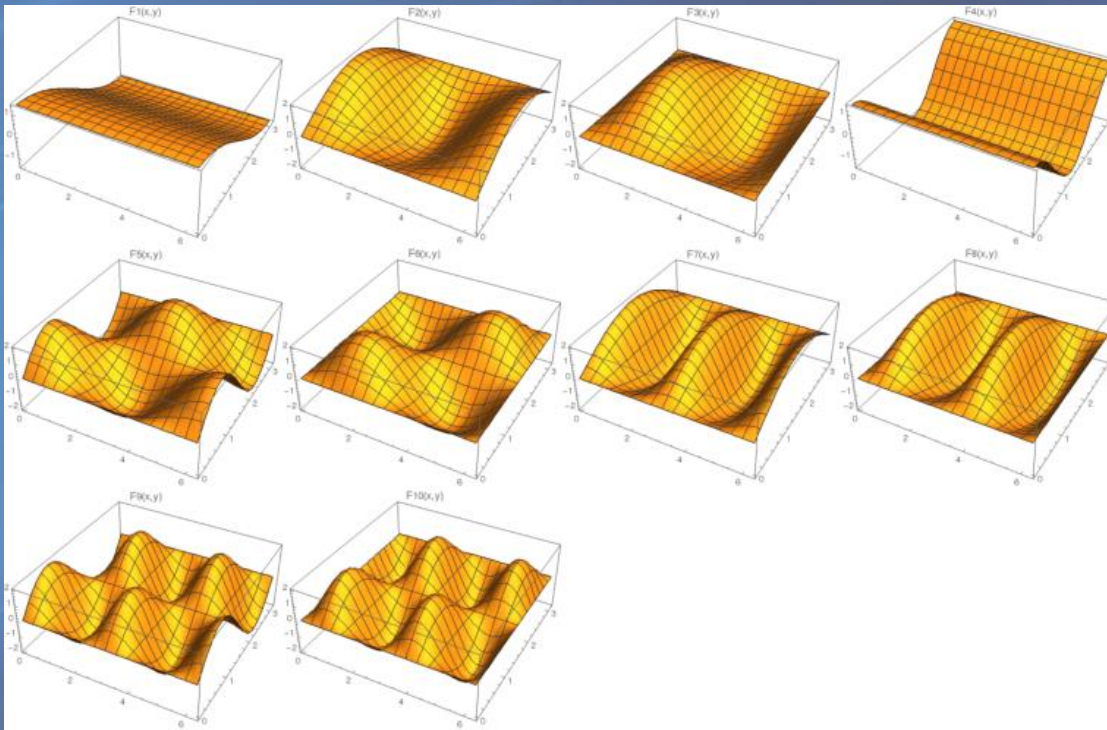
36-variable model



New version of the model: MAOSOAM

Rectangular geometry

Channel flow for the ocean too



Low-order system of 40 variables, 20 for the ocean and 20 for the atmosphere

Parameter values

Table 1. Dimensional Parameters Present in the Coupled Ocean-Atmosphere Models^a

Dynamic Atmosphere	Dynamic Ocean	Geometry	Coupling
$k_d = gC/\Delta p \text{ s}^{-1}$	$L_R = (g'H)^{1/2}/f_0 n$	$L_y = \pi L = 5000 \text{ km}$	$\epsilon_a = 0.7$
$k'_d = k_d$	$r = 10^{-7} \text{ s}^{-1}$	$n = (2L_y)/L_x = 1.5$	$\lambda = c_{p,a} C \text{ W m}^{-2} \text{ K}^{-1}$
$\sigma = 2.16 \cdot 10^{-6} \text{ J kg}^{-1} \text{ Pa}^{-2}$	$\gamma_o = c_{p,o} \rho_o H \text{ J m}^{-2} \text{ K}^{-1}$	$f_0 = 0.0001052 \text{ s}^{-1}$	$d = C/(\rho_o H) \text{ s}^{-1}$
$\gamma_a = 10^7 \text{ J m}^{-2} \text{ K}^{-1}$		$\beta = 1.62 \cdot 10^{-11} \text{ m}^{-1} \text{ s}^{-1}$	
		$\Delta p = 500 \text{ hPa}$	

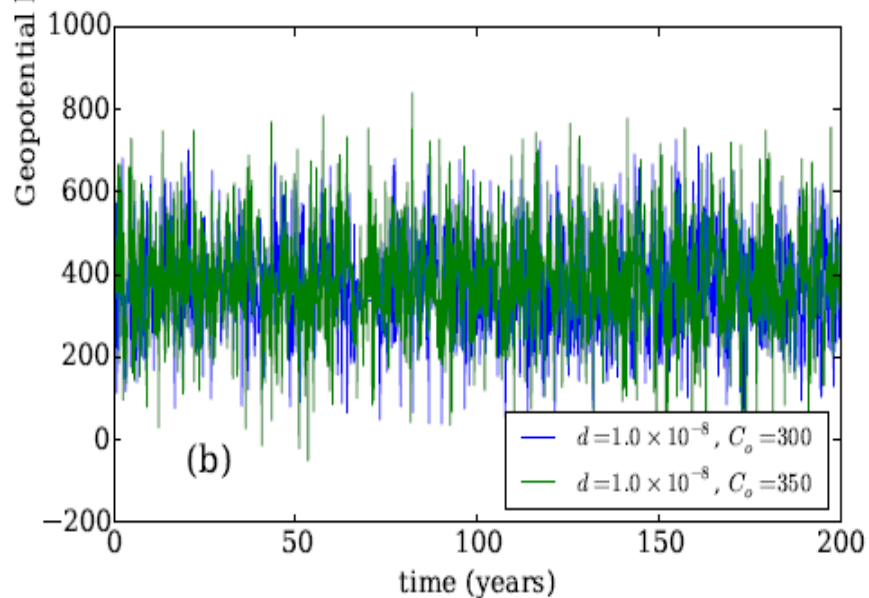
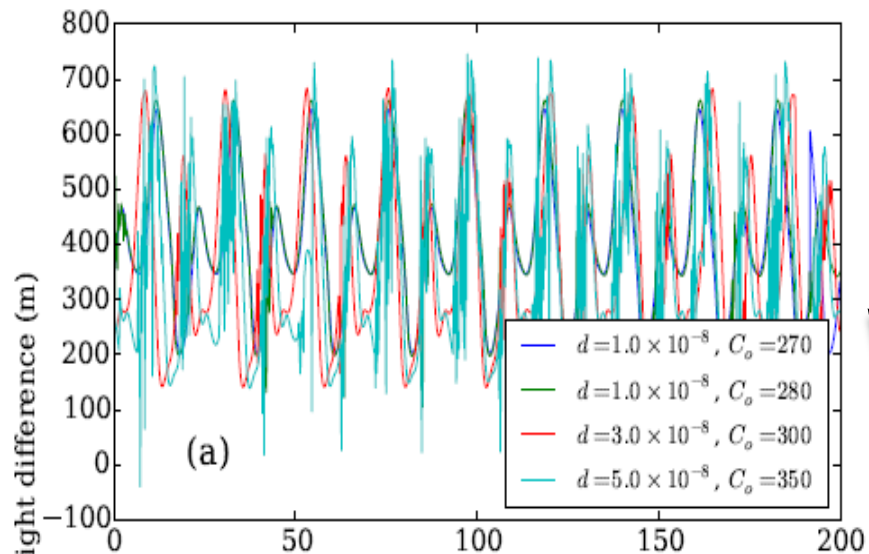
^a $c_{p,a}$ and σ_B are the usual specific heat at constant pressure of the air and the Stefan-Boltzmann constant, fixed to $1004 \text{ J kg}^{-1} \text{ K}^{-1}$ and $5.6 \cdot 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$, respectively. The density, ρ_o , and the specific heat at constant pressure, $c_{p,o}$, for the ocean layer are fixed to 1000 kg m^{-3} and $4000 \text{ J kg}^{-1} \text{ K}^{-1}$. g and g' are the gravity and reduced gravity fixed to 10 and 0.031 m s^{-2} , respectively.

Vannitsem, 2015, Geophys Res Lett

4 important parameters: n , C , H and $C_0 \equiv S_0$

Variability and Lyapunov instability properties of the coupled ocean-atmosphere system

Solution of the VDDG model

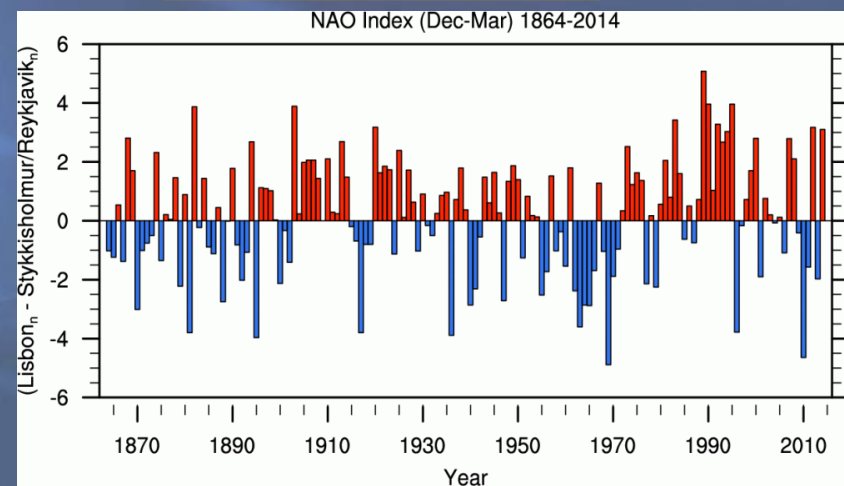


Geop. Height. Diff. between
two points from the North
and South parts of the domain

Vannitsem et al, 2015, Physica D

$n=1.5$
 $H=500$ m

Decadal Variability



Bifurcation diagram

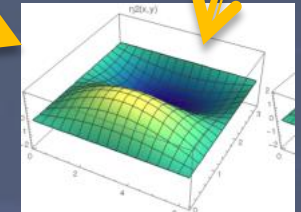
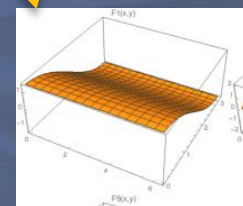
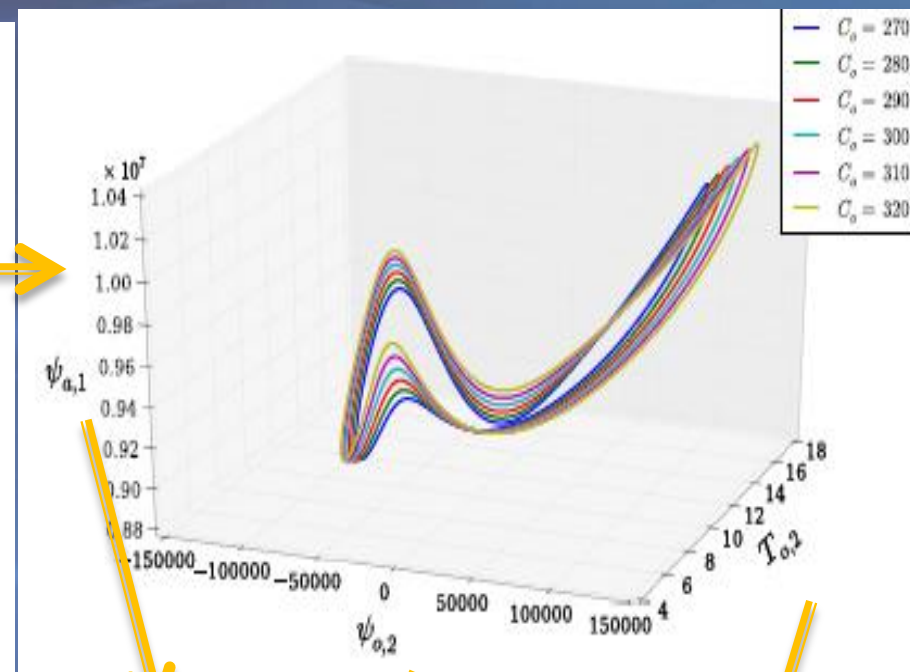
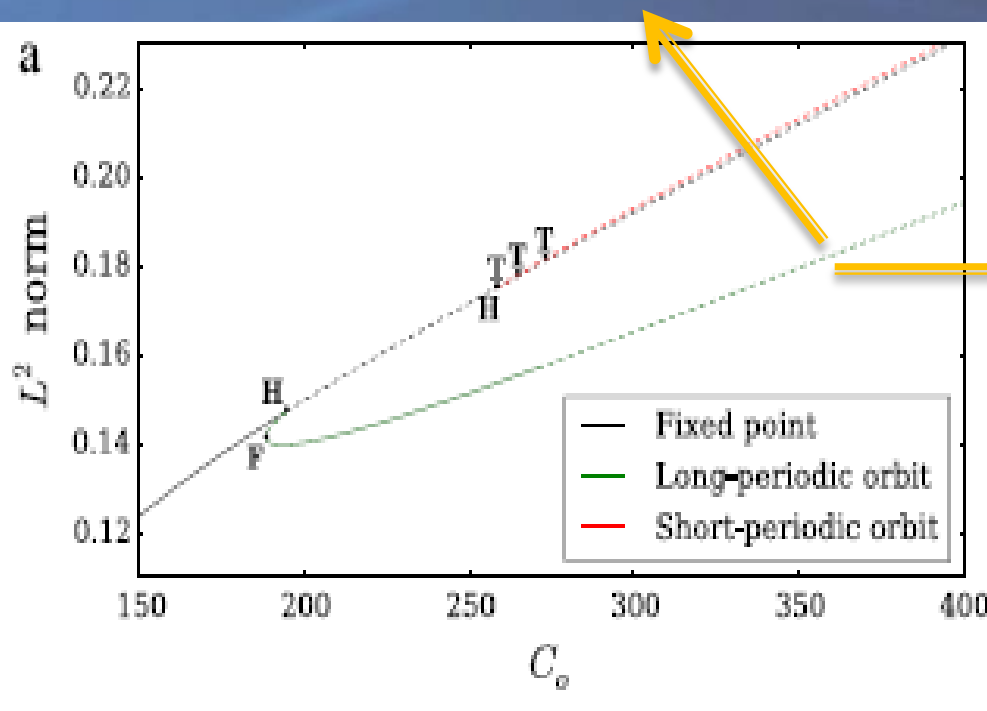
$n=1.5$

$H=500 \text{ m}$

$\lambda=20 \text{ W m}^{-2} \text{ K}^{-1}$

$d=10^{-8} \text{ s}^{-1}$

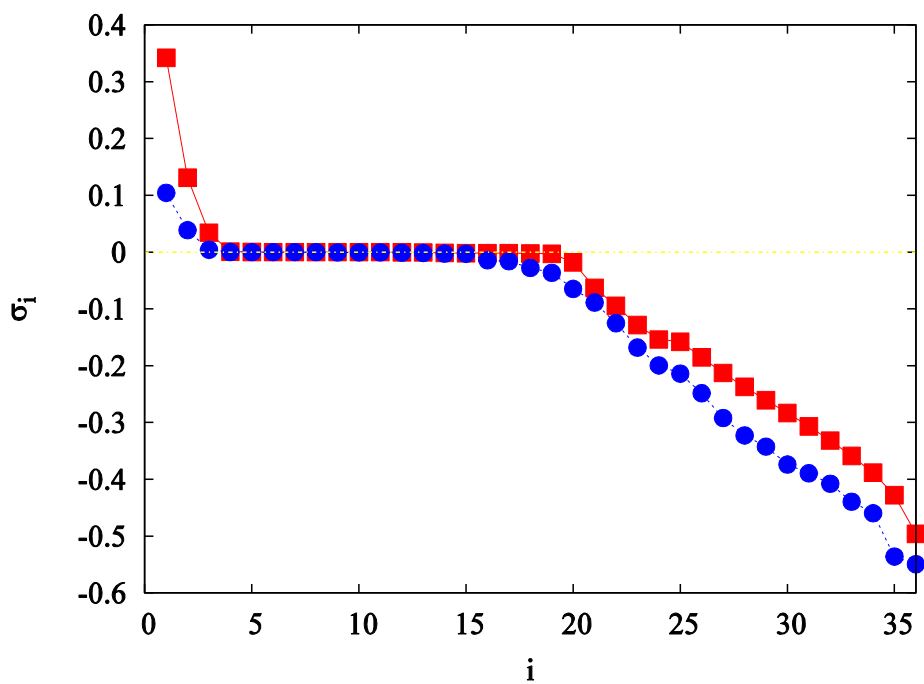
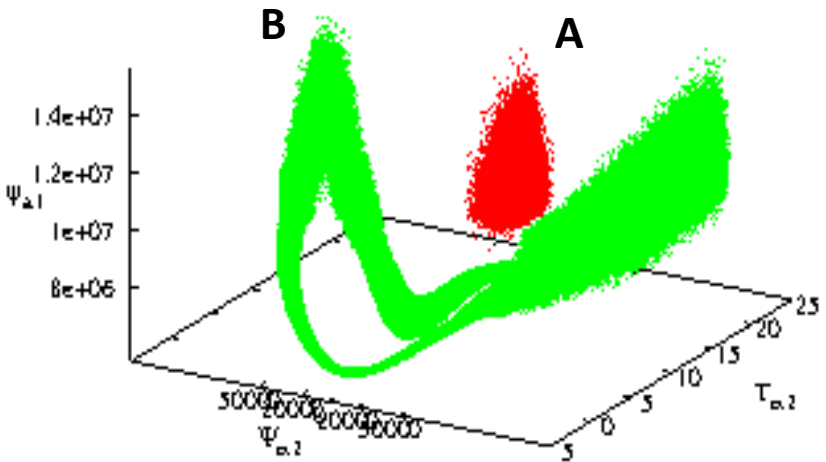
Slow branch



Consider 2 different attractors
(obtained with different parameters)

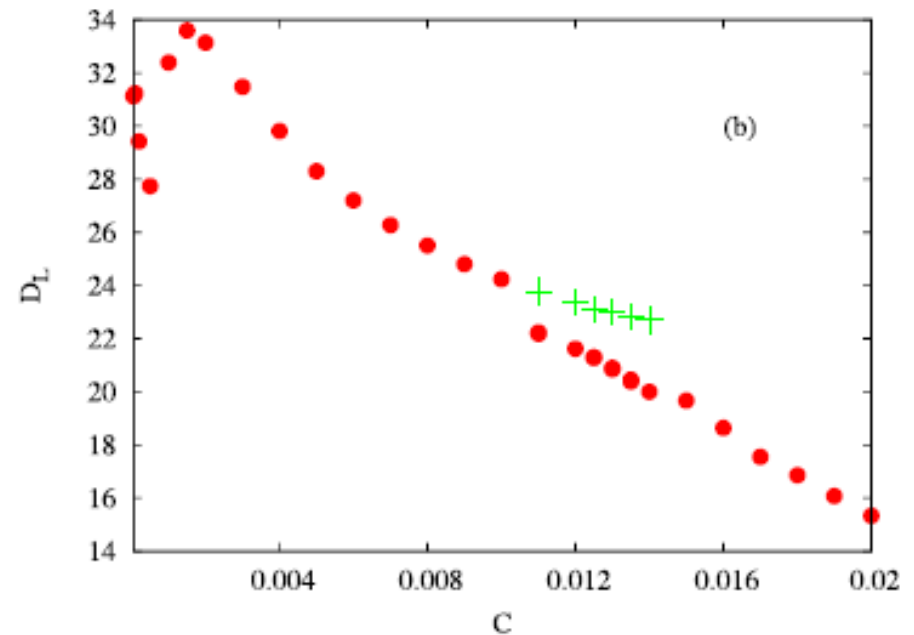
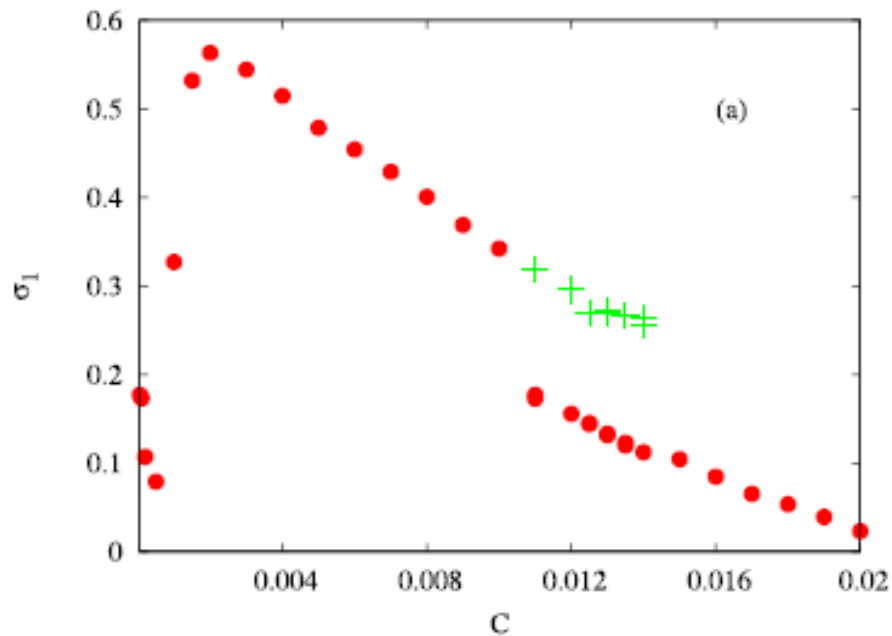
$n=1.5$
 $H=164 \text{ m}$
 $C_0=S_0=310 \text{ W/m}^2$

Lyapunov spectra



Vannitsem, 2017, Chaos, 27, 032101

Changes of the Lyapunov instability as a function of friction



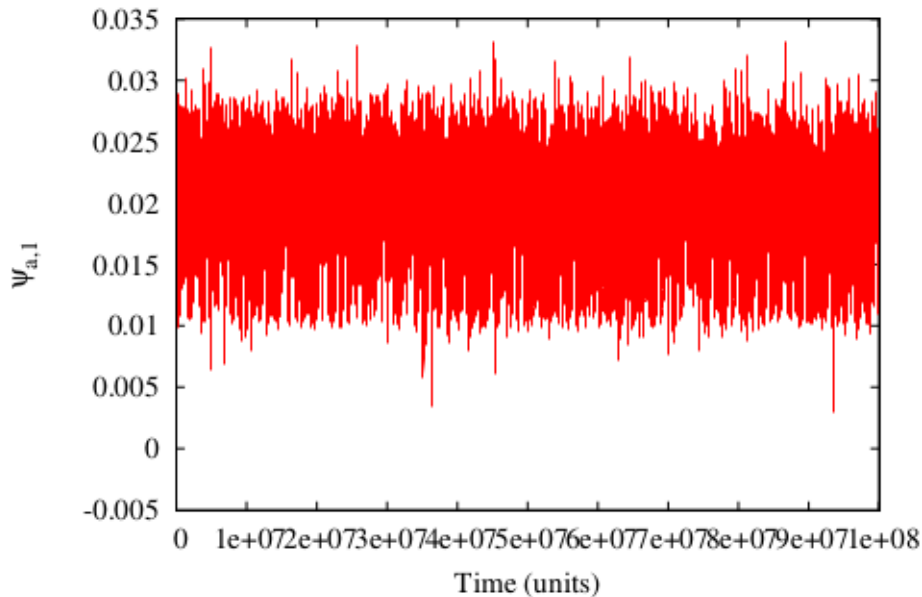
Vannitsem, 2017, Chaos, 27, 032101

$n=1.5$
 $H=164 \text{ m}$
 $C_0=S_0=310 \text{ W/m}^2$

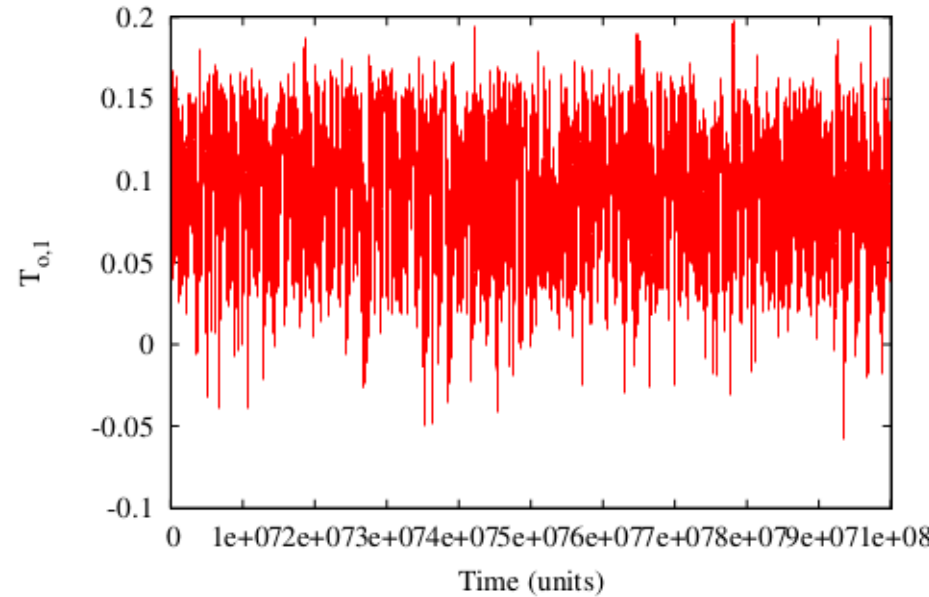
Results with the new version of the model, MAOSOAM

Depth of the ocean= 1000 m, $S_o = 350 \text{ W/m}^2$ $n = 2 L_y/L_x = 1.5$

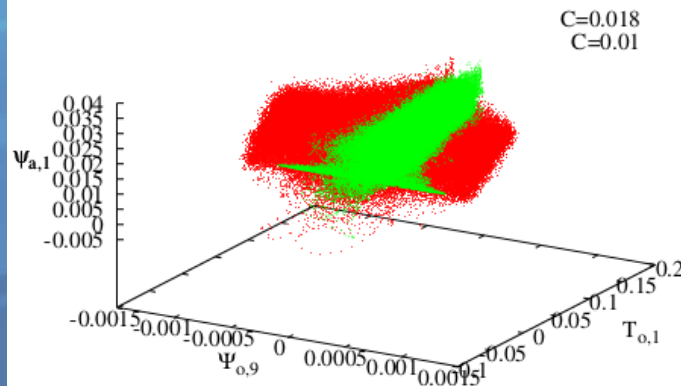
$N=1.5$, $C=0.01$, $H=1000$, $C_o = 350$, No seasonality



$N=1.5$, $C=0.01$, $H=1000$, $C_o = 350$, No seasonality



$N=1.5$, $C=0.01$ and $C=0.018$, $H=1000$, $C_o = 350$, No seasonality

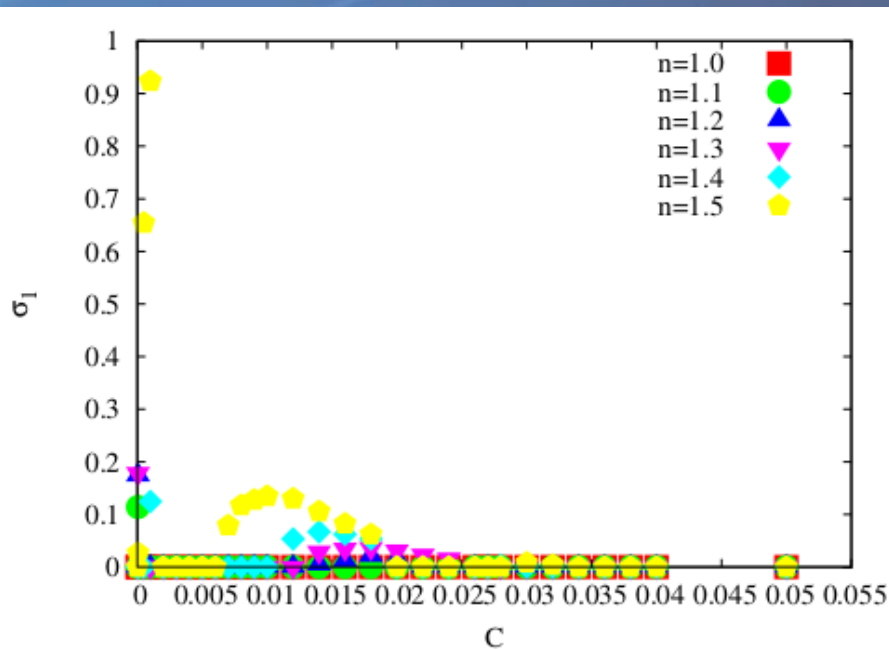


Results with the new version of the model, MAOSOAM

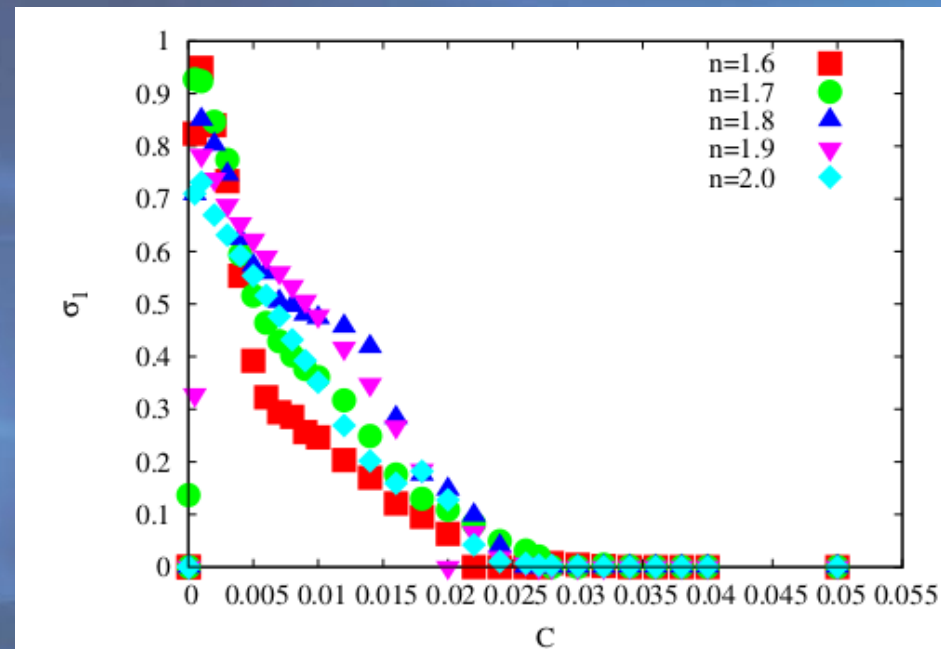
Dominant Lyapunov exponent for different values of C , for different aspect ratios n ,

$$n = 2 L_y / L_x$$

Depth of the ocean= 1000 m, $S_o = 350 \text{ W/m}^2$



The interaction with the ocean induces chaos



The interaction with the ocean reduces chaos

The background of the slide is a deep blue color with several lighter blue, wavy, translucent lines that sweep across the frame from the bottom left towards the top right, creating a sense of fluid motion or atmospheric currents.

Error dynamics of atmospheric and climate flows

Typical evolution of the averaged Error

A lot of discussions on the shape of the mean square error evolution

$$\langle E_t^2 \rangle = \int d\mathbf{x}_0 \rho(\mathbf{x}_0) \{y(t; \mathbf{x}_0 + \boldsymbol{\varepsilon}) - \mathbf{x}(t; \mathbf{x}_0)\}^2$$

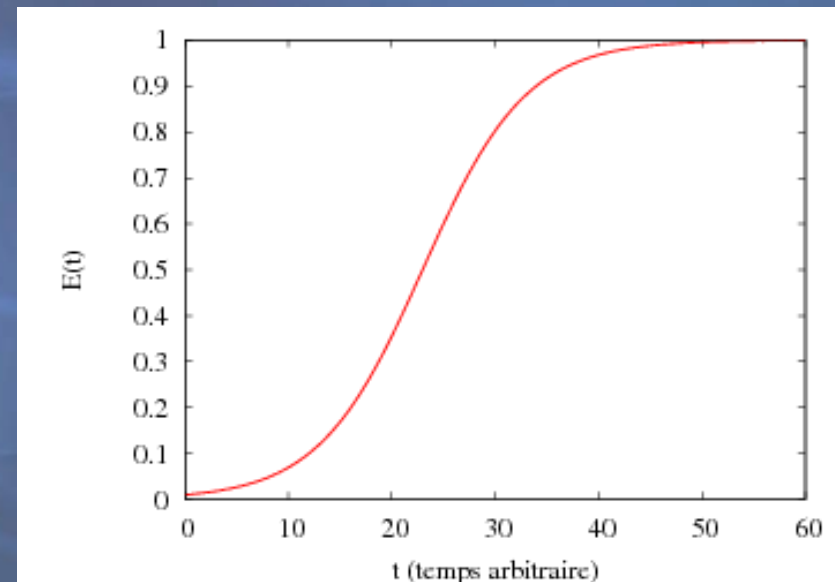
Idealized shape of the evolution of the error :

$$\frac{dE}{dt} = aE - bE^2$$

Lorenz (1969, Tellus, 21)

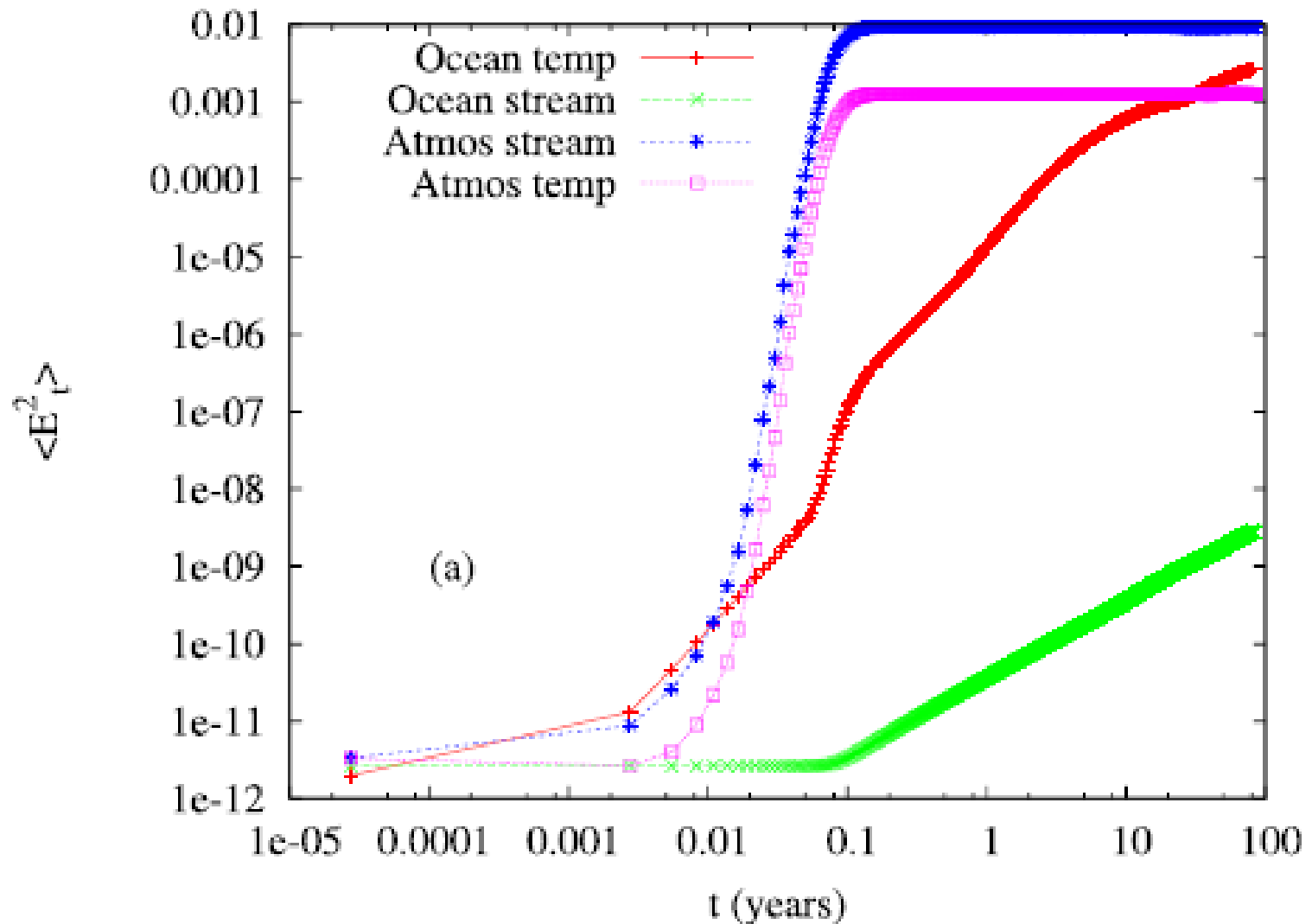
Trevisan et al (1992, JAS, 49)

Nicolis (1992, QJRMS, 118)



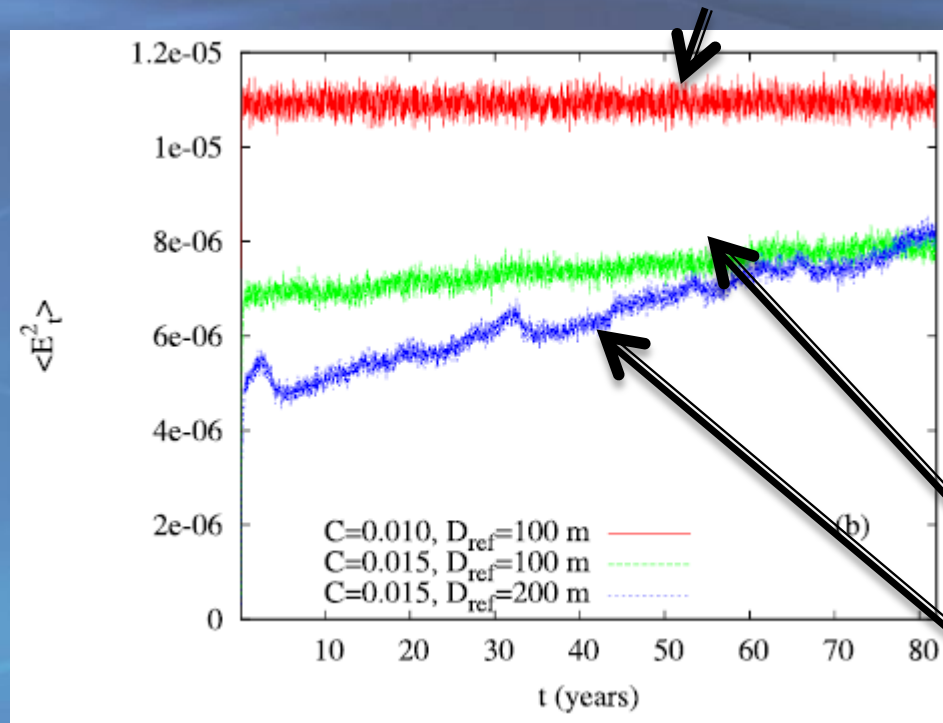
Error dynamics in the VDDG model

Attractor A



Error dynamics in VDDG model

No LFV



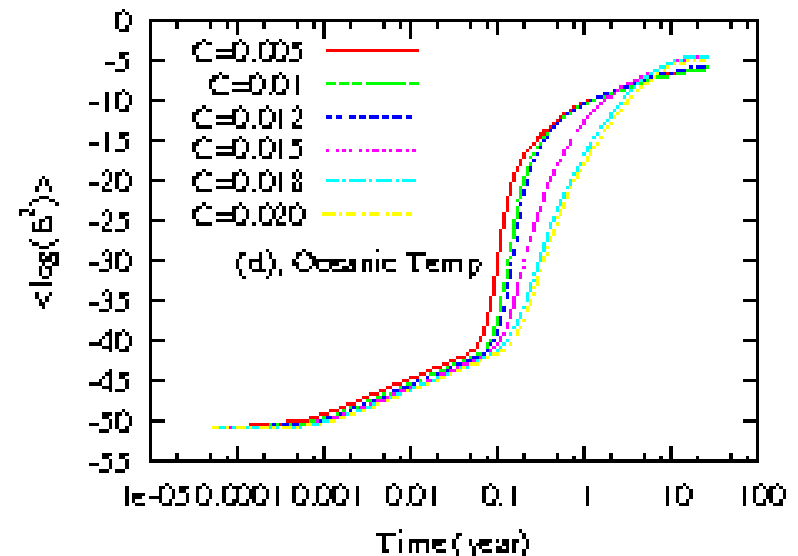
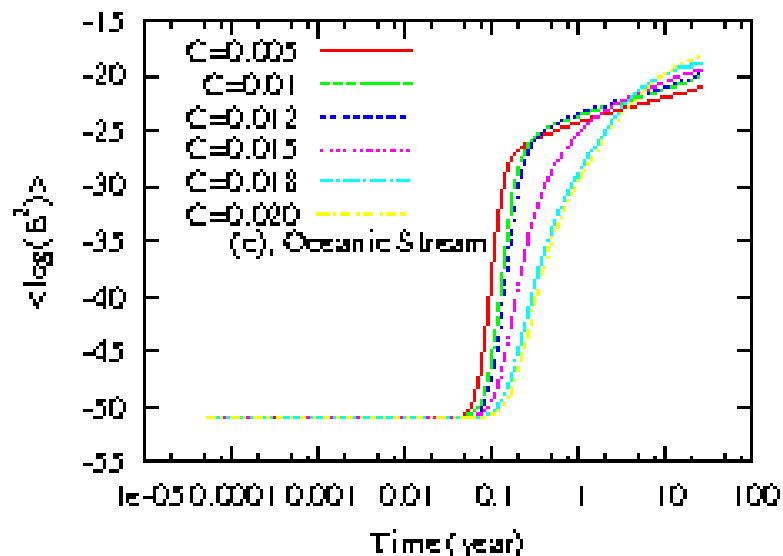
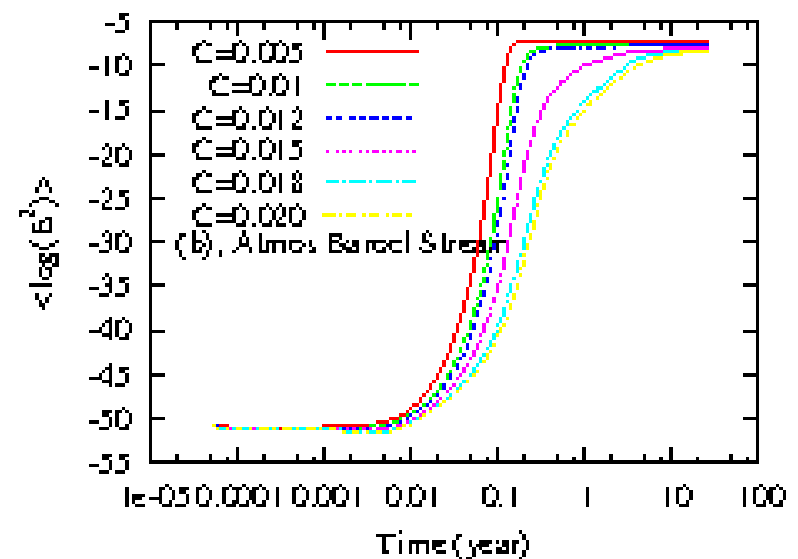
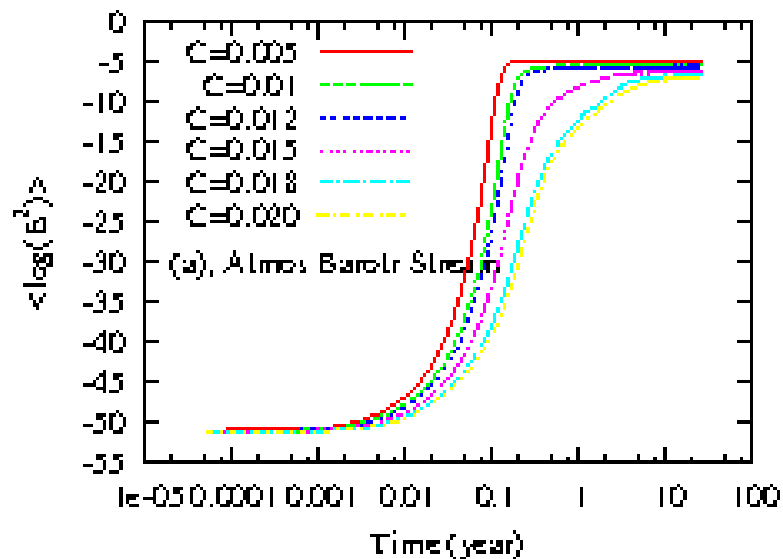
$\psi_{a,1}$

With LFV

Results with the new version of the model, MAOSOAM

Depth of the ocean= 1000 m, $S_0 = 350 \text{ W/m}^2$

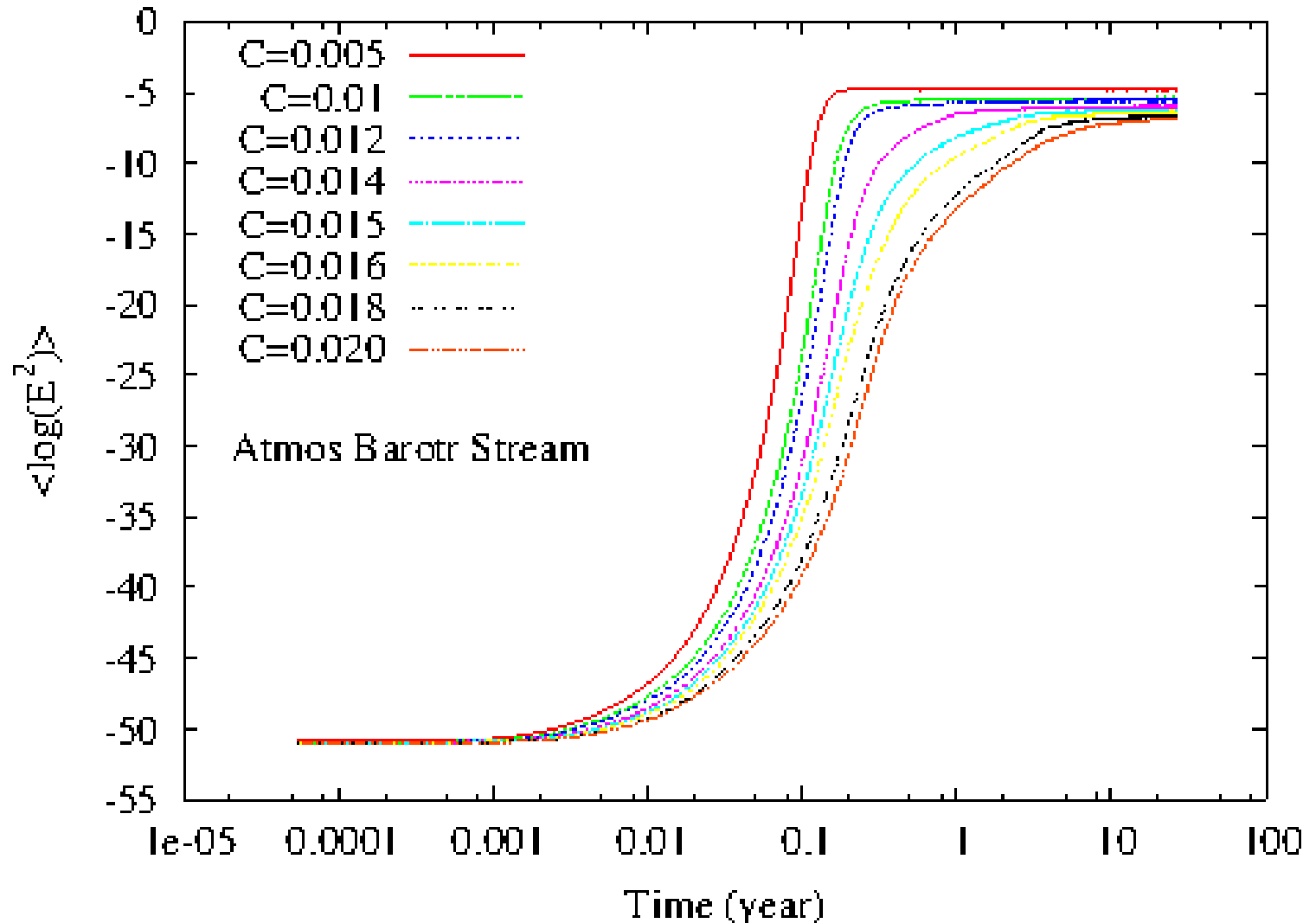
$$n = 2 L_y / L_x = 1.7$$



Results with the new version of the model, MAOSOAM

Depth of the ocean= 1000 m, $S_o = 350 \text{ W/m}^2$

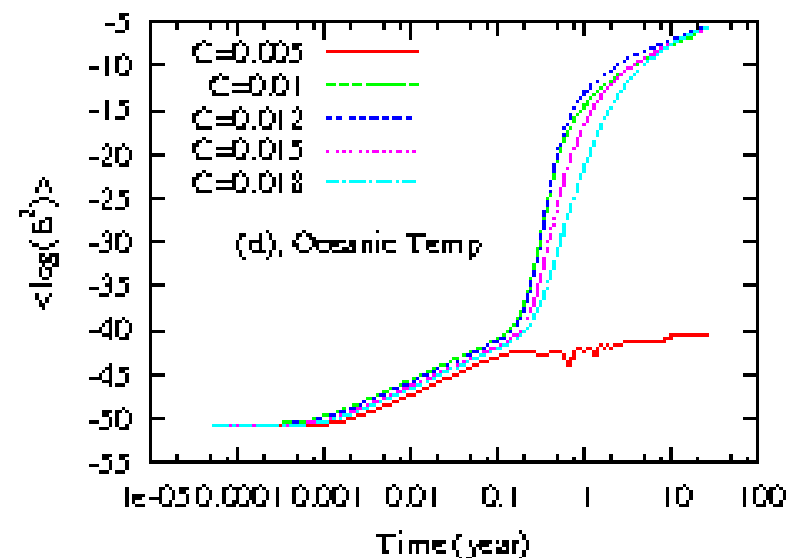
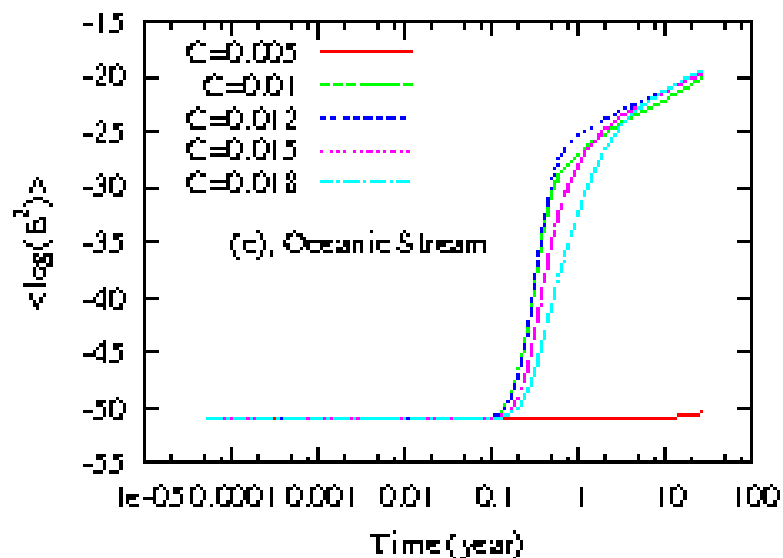
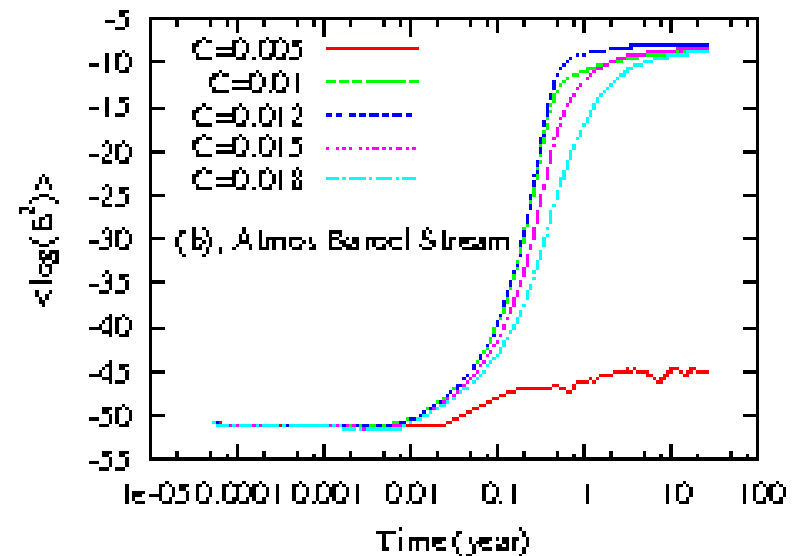
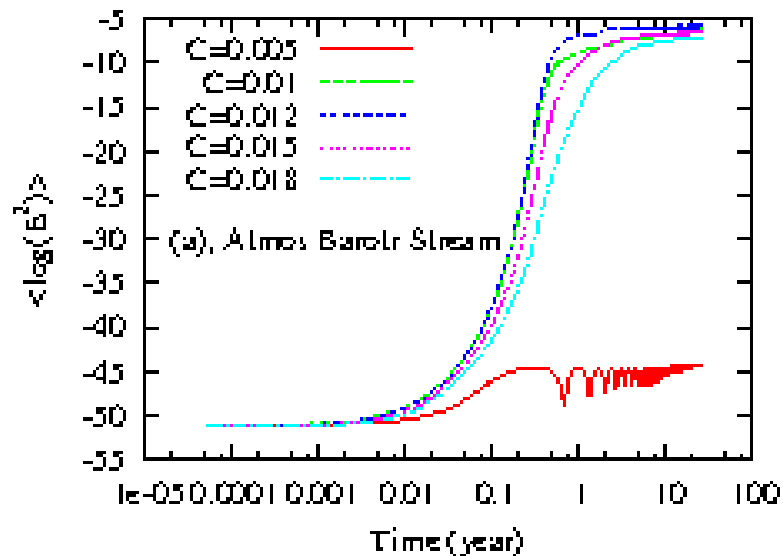
$$n = 2 L_y / L_x = 1.7$$



Results with the new version of the model, MAOSOAM

Depth of the ocean= 1000 m, $S_0 = 350 \text{ W/m}^2$

$$n = 2 L_y / L_x = 1.5$$



Some conclusions

In the model version with closed boundaries in the ocean (VDDG or MAOOAM, $n=1.5$):

- Friction is tempering chaos.
- The long term predictability of the atmosphere is related to the strong coupling between the ocean and the atmosphere (in the Low-order O-A model), i.e. a coupled mode should be present

In the new model version with channel flow:

- The influence of friction depends on the aspect ratio, n , with a chaos-induced dynamics for small aspect ratio and a chaos-tempered dynamics for large aspect ratio
- The long term predictability of the atmosphere depends on the proximity to bifurcation points? (Still work to be done to clarify this feature)

Future investigations

- Analysis of a higher order coupled ocean-atmosphere system
- Coupling with other components of the climate system

Some references

De Cruz, L., J. Demaeyer and S. Vannitsem, **A modular arbitrary-order ocean-atmosphere model: MAOOAM V1.0**, Geoscientific Model Development, 9, 2793-2808 , 2016. (GITHUB)

De Cruz, L., Schubert, S., Demaeyer, J., Lucarini, V., and Vannitsem, S.: **Exploring the Lyapunov instability properties of high-dimensional atmospheric and climate models**, Nonlin. Processes Geophys., <https://doi.org/10.5194/npg-2017-76>, in press, 2018.

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Vannitsem S. and V. Lucarini, **Statistical and Dynamical Properties of Covariant Lyapunov Vectors in a Coupled Atmosphere-Ocean Model - Multiscale Effects, Geometric Degeneracy, and Error Dynamics**, J. Phys. A., 2016. ArXiv: [arXiv:1510.00298v3](https://arxiv.org/abs/1510.00298v3)